

Shi(f)t happens

- Krylov methods for shifted linear systems -

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SIAM Student Krylov Day 2015

What's a shifted linear system?

Definition

Shifted linear systems are of the form

$$(A - \omega_k I) \mathbf{x}_k = \mathbf{b},$$

where $\{\omega_k\}_{k=1}^N \in \mathbb{C}$ are a sequence of (many) *shifts*.

For the simultaneous solution, **Krylov methods** are well-suited because of the *shift-invariance* property:

$$\mathcal{K}_m(A, \mathbf{b}) \equiv \text{span}\{\mathbf{b}, A\mathbf{b}, \dots, A^{m-1}\mathbf{b}\} = \mathcal{K}_m(A - \omega I, \mathbf{b}).$$

"Proof" (shift-invariance)

For $m = 2$: $\mathcal{K}_2(A, \mathbf{b}) = \text{span}\{\mathbf{b}, A\mathbf{b}\}$

$$\mathcal{K}_2(A - \omega I, \mathbf{b}) = \text{span}\{\mathbf{b}, A\mathbf{b} - \omega\mathbf{b}\} = \text{span}\{\mathbf{b}, A\mathbf{b}\}$$

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Some (semi) open research questions...

$$(A - \omega_k I) \mathbf{x}_k = \mathbf{b}$$

- 1 Can we allow multiple right-hand sides?
- 2 Which preconditioners preserve shift-invariance?
- 3 Can we apply restarting and nested algorithms?
- 4 Can we benefit from (spectral) deflation?
- 5 What if $(A - \omega_k M) \mathbf{x}_k = \mathbf{b}$?
- 6 Where do shifted systems occur in practice?

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Outline

- 1 Multi-shift Krylov methods
- 2 Some words on preconditioning
- 3 Nested multi-shift Krylov methods
- 4 Geophysical applications

Multi-shift GMRES

After m steps of Arnoldi, we have,

$$AV_m = V_{m+1}\underline{H}_m,$$

and the approximate solution yields:

$$\mathbf{x}_m \approx V_m \mathbf{y}_m, \quad \text{where } \mathbf{y}_m = \underset{\mathbf{y} \in \mathbb{C}^m}{\operatorname{argmin}} \|\underline{H}_m \mathbf{y} - \|\mathbf{b}\| \mathbf{e}_1\|.$$

For **shifted systems**, we get

$$(A - \omega I)V_m = V_{m+1}(\underline{H}_m - \omega \underline{I}_m),$$

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Preconditioning is a problem

Main disadvantage:

Preconditioners are in general not easy to apply. For

$$(A - \omega I)P_{\omega}^{-1}y^{(\omega)} = b, \quad P_{\omega}x^{(\omega)} = y^{(\omega)}$$

it does **not** hold:

$$\mathcal{K}_m(AP^{-1}, b) \neq \mathcal{K}_m(AP_{\omega}^{-1} - \omega P_{\omega}^{-1}, b).$$

However, there are ways...

Reference

B. Jegerlehner, *Krylov space solvers for shifted linear systems*. Published online arXiv:hep-lat/9612014, 1996.

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Preconditioning is a problem

... or has been a problem ?

Short historical overview:

2002 Shift-and-invert preconditioner:

$$\mathcal{P} = (A - \tau I), \quad \tau \approx \{\omega_1, \dots, \omega_N\}$$

2007 Many shift-and-invert preconditioners:

$$\mathcal{P}_j = (A - \tau_j I)$$

2013 Polynomial preconditioners:

$$p_n(A) \approx A^{-1}, \quad p_n^\omega(A) \approx (A - \omega I)^{-1}$$

2014 Nested Krylov methods

Nested multi-shift Krylov methods

Methodology:

- Martin knows: Polynomial preconditioners exist
- Question: Can we use a Krylov polynomial?

Nested multi-shift Krylov methods:

- Use an inner multi-shift Krylov method as preconditioner.
- For inner method, require collinear residuals $[\mathbf{r}_j^{(\omega)} = \gamma \mathbf{r}_j]$.

This is the case for:

- ▶ multi-shift GMRES [1998]
 - ▶ multi-shift FOM [2003]
 - ▶ multi-shift BiCG [2003]
 - ▶ multi-shift IDR(s) [new!]
- Using γ , we can preserve the shift-invariance in the outer Krylov iteration.

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Nested multi-shift Krylov methods:

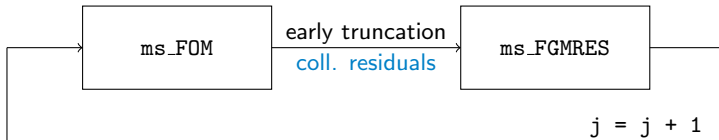
- Use an **inner multi-shift Krylov method** as preconditioner.
- For inner method, require collinear residuals $[\mathbf{r}_j^{(\omega)} = \gamma \mathbf{r}_j]$.

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Nested multi-shift Krylov methods

Inner-outer iteration:



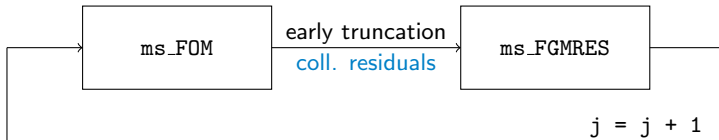
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Multi-shift FOM as inner method

Classical result: In FOM, the residuals are

$$\mathbf{r}_j = \mathbf{b} - A\mathbf{x}_j = \dots = -h_{j+1,j}\mathbf{e}_j^T \mathbf{y}_j \mathbf{v}_{j+1}.$$

Thus, for the shifted residuals it holds:

$$\mathbf{r}_j^{(\omega)} = \mathbf{b} - (A - \omega I)\mathbf{x}_j^{(\omega)} = \dots = -h_{j+1,j}^{(\omega)}\mathbf{e}_j^T \mathbf{y}_j^{(\omega)} \mathbf{v}_{j+1},$$

which gives $\gamma = y_j^{(\omega)} / y_j$.

Reference

V. Simoncini, *Restarted full orthogonalization method for shifted linear systems*. BIT Numerical Mathematics, 43 (2003).

Flexible multi-shift GMRES as outer method

Use flexible GMRES in the outer loop,

$$(A - \omega I) \hat{V}_m = V_{m+1} \underline{H}_m^{(\omega)},$$

where one column yields

$$(A - \omega I) \underbrace{\mathcal{P}(\omega)_j^{-1} \mathbf{v}_j}_{\text{inner loop}} = V_{m+1} \underline{\mathbf{h}}_j^{(\omega)}, \quad 1 \leq j \leq m.$$

The “inner loop” is the truncated solution of $(A - \omega I)$ with right-hand side $\mathbf{v}_j$ using [msFOM](#).

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Flexible multi-shift GMRES

The inner residuals are:

$$\begin{aligned}\mathbf{r}_j^{(\omega)} &= \mathbf{v}_j - (A - \omega I)\mathcal{P}(\omega)_j^{-1}\mathbf{v}_j, \\ \mathbf{r}_j &= \mathbf{v}_j - A\mathcal{P}_j^{-1}\mathbf{v}_j,\end{aligned}$$

Imposing $\mathbf{r}_j^{(\omega)} = \gamma \mathbf{r}_j$ yields:

$$(A - \omega I)\mathcal{P}(\omega)_j^{-1}\mathbf{v}_j = \gamma A\mathcal{P}_j^{-1}\mathbf{v}_j - (\gamma - 1)\mathbf{v}_j \quad (*)$$

Note that the right-hand side in $(*)$ is a preconditioned shifted system!

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which yields:

$$\underline{H}_m^{(\omega)} = (\underline{H}_m - \underline{I}_m) \underline{\Gamma}_m + \underline{I}_m,$$

with $\underline{\Gamma}_m := \text{diag}(\gamma_1, \dots, \gamma_m)$.

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Geophysical applications

The model

The time-harmonic elastic wave equation

For **many** (angular) frequencies ω_k , we solve

$$-\omega_k^2 \rho(\mathbf{x}) \hat{\mathbf{u}} - \nabla \cdot \sigma(\hat{\mathbf{u}}, c_p, c_s) = \hat{\mathbf{s}}, \quad \mathbf{x} \in \Omega \subset \mathbb{R}^{2,3},$$

together with absorbing or reflecting boundary conditions.

Inverse (discrete) Fourier transform:

$$\mathbf{u}(\mathbf{x}, t) = \sum_k \hat{\mathbf{u}}(\mathbf{x}, \omega_k) e^{i\omega_k t}$$

The **discretized** time-harmonic elastic wave equation is quadratic in ω_k :

$$(K + i\omega_k C - \omega_k^2 M)\underline{\hat{\mathbf{u}}} = \underline{\hat{\mathbf{s}}},$$

which can be re-arranged as,

$$\left[\begin{pmatrix} iM^{-1}C & M^{-1}K \\ I & 0 \end{pmatrix} - \omega_k \begin{pmatrix} I & 0 \\ 0 & I \end{pmatrix} \right] \begin{pmatrix} \omega_k \underline{\hat{\mathbf{u}}} \\ \underline{\hat{\mathbf{u}}} \end{pmatrix} = \begin{pmatrix} M^{-1}\underline{\hat{\mathbf{s}}} \\ 0 \end{pmatrix}.$$

The latter is of the form:

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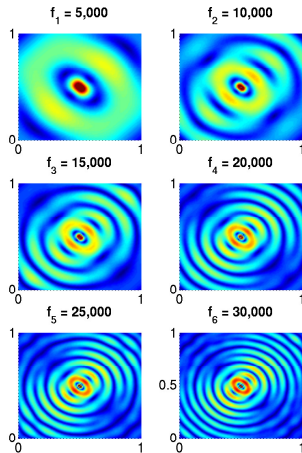
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A first example - The setting

Test case from literature:

- $\Omega = [0, 1] \times [0, 1]$
- $h = 0.01$ implying
 $n = 10.201$ grid points
- system size:
 $4n = 40.804$
- $N = 6$ frequencies
- point source at center

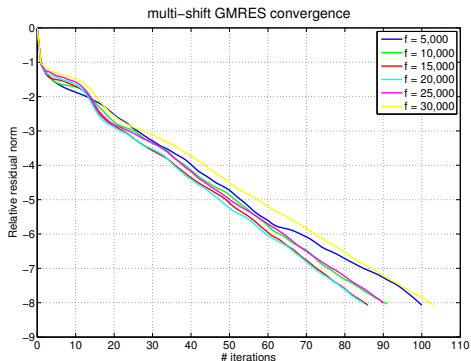


Reference

- T. Airaksinen, A. Pennanen, and J. Toivanen, *A damping preconditioner for time-harmonic wave equations in fluid and elastic material*. Journal of Computational Physics, 2009.

A first example - Convergence behavior (1/2)

Preconditioned **multi-shift GMRES**:

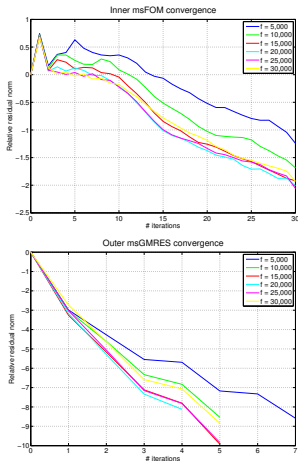


We observe:

- simultaneous solve
- CPU time: 17.71s

A first example - Convergence behavior (2/2)

Preconditioned **nested FOM-FGMRES**:



We observe:

- 30 inner iterations
- truncate when inner residual norm ~ 0.1
- very few outer iterations
- CPU time: **9.62s**

Summary

- ✓ Nested Krylov methods for $A\mathbf{x} = \mathbf{b}$ are widely used
↪ extension to shifted linear systems is possible
- ✓ Multiple combinations of inner-outer methods possible, e.g. FOM-FGMRES, IDR-FQMRIDR, ...
- ✓ The shift-and-invert preconditioner (or the polynomial preconditioner) can be applied on top
- ? At the moment: Multiple right-hand sides (with K. Soodhalter, U Linz)

Thank you for your attention!